

Probing novel TeV physics through precision calculations of scalar and tensor charges of the nucleon

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We present an update on the calculation of matrix elements of iso-vector scalar, axial and tensor charges between a neutron and a proton state. These matrix elements are needed to probe novel scalar and tensor interactions in neutron beta-decay that can arise in extensions of the Standard Model at the TeV scale. Our calculations are being done using valence clover fermions on dynamical $N_f = 2 + 1 + 1$ HISQ configurations generated by the MILC Collaboration. We provide preliminary estimates of the dependence of these matrix elements on the light quark masses, lattice spacing, and the time separation between the source and sink of the nucleons. We also find that the renormalization constants calculated using the RI-sMOM scheme are close to unity for the HYP smeared HISQ lattices.

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1. Introduction

The observed electroweak symmetry breaking in the Standard Model (SM) points to new physics at the TeV scale, which is currently being explored at the LHC. There are many candidate extensions of the SM, such as supersymmetry and extra dimensions, but, so far, there is no experimental guidance on what new interactions and particles exist at this scale. A theoretical approach (effective field theory) is to postulate new interactions at the TeV scale (without necessarily specifying a model) and analyze the corrections they would give rise to in low-energy (few-GeV) physics. Precision measurements of decays can then be used to bound the allowed parameter space of the posulated interactions. In Ref. [1], we have shown that precision measurements of decays of (ultra)cold neutrons at the 10^{-3} level can be used to improve the bounds on new scalar and tensor interactions at the TeV scale provided estimates accurate to 10–20% of the matrix elements of isovector scalar and tensor bilinear quark operators are realized.

In these proceedings, we will summarize the progress we have made in the needed calculations to extract the charges g_A , g_S and g_T . These charges are defined as

$$g_A = Z_A \langle p | O_A | n \rangle ; \quad g_S = Z_S \langle p | O_S | n \rangle ; \quad g_T = Z_T \langle p | O_T | n \rangle ,$$

where Z_Γ are the corresponding renormalization constants, which we calculate nonperturbatively in the RI-sMOM scheme [2, 3]. This report is a follow up to Ref. [1], and we will use the notation established there. The calculations are being done using the 2+1+1 flavor HISQ lattices being generated by the MILC Collaboration [4]. Clover fermions on HYP smeared HISQ lattices are being used for constructing the correlation functions. Three values of the lattice spacing are being analyzed, $a = 0.12, 0.09$ and 0.06 fm, to perform the continuum extrapolation, and at each lattice spacing we analyze two values of the light quark masses corresponding to $M_\pi \approx 310$ and 220 MeV. We will address the following sources of statistical and systematic errors in the extraction of matrix elements and, from them, the charges.

- The signal and the statistical errors in the two- and three-point correlation functions.
- The dependence of $g_{A,S,T}$ on the light-quark (pion) masses and lattice spacing to estimate uncertainty due to chiral and continuum extrapolations.
- The effect of contamination by excited states.
- Estimates of the renormalization constants calculated in the RI-sMOM scheme.

A similar analysis has been reported in Ref. [5] at this conference.

2. Statistics

The MILC Collaboration has produced ensembles of roughly 5500 trajectories of 2+1+1 flavor HISQ lattices at each of the six values of quark masses and lattice spacings we are analyzing as described in Table 1. Five hundred trajectories are discarded for thermalization. Configurations are then analyzed separated by five trajectories. On each configuration, four smeared sources, displaced both in time and space directions to reduce correlations, are used. Furthermore, two sets

of these four source points, again maximally separated in space and time directions, are used on each alternate configuration to reduce correlations. The roughly 500 configurations with each of these two sets of sources are also analyzed separately. We verify that the two sets give compatible results and the errors are roughly $\sqrt{2}$ larger compared to the full set. Our main conclusion is that with 1000 or more configurations one gets a statistically significant signal needed for obtaining the design accuracy of 10–20% in both the 2-point and 3-point functions for all cases. Extraction of g_S is the nosiest and drives the overall sensitivity as discussed below.

a (fm)	m_l/m_s	M_π (MeV)	$L^3 \times T$	$M_\pi L$	Confs (Z)	Confs (ME)	Δt
0.12	0.2	305	$24^3 \times 64$	4.54	50	1013/1013	8,9,10,11,12
0.12	0.1	217	$32^3 \times 64$	4.29	50	958/958	8,10,12
0.09	0.2	313	$32^3 \times 96$	4.5		391/1000	12,14
0.09	0.1	220	$48^3 \times 96$	4.73		443/1000	12,14
0.06	0.2	320	$48^3 \times 144$	4.53		330/1000	16,20
0.06	0.1	229	$64^3 \times 144$	4.28		0/1000	

Table 1: Lattice parameters of the six ensembles of 2+1+1 flavor HISQ lattices analyzed. We give a status report (as of November 2012) on the number of analyzed and total configurations used in calculations of ME and for the calculation of renormalization constants (Z). The last column gives in lattice units the time separations Δt between source and sink investigated to understand the excited-state contamination.

3. Excited-State Contamination

The desired matrix elements need to be calculated between ground state nucleons. The operators used to create and annihilate the states, however, couple to the nucleon and all its excited states. There are two possible ways to reduce contribution from excited states: by reducing the overlap of the interpolating operator with the excited states and by increasing the time separation Δt between the source and sink to exponentially suppress excited-state contamination. Statistics limit the upper value of Δt that can be explored. We use smeared sources and sinks to reduce overlap and investigate up to five time separations.

Assuming that all excited-state contamination can be represented by a single excited state with amplitude A_1 and mass M_1 , we can write the 3-point function with source at $t = 0$, operator insertion at $t = t$ and sink at $t = \Delta t$ as

$$\begin{aligned} \mathcal{C}_\Gamma = & |A_0|^2 \langle 0|O_\Gamma|0\rangle e^{-M_0\Delta t} + |A_1|^2 \langle 1|O_\Gamma|1\rangle e^{-M_1\Delta t} + \\ & A_0 A_1^* \langle 0|O_\Gamma|1\rangle e^{-M_0 t} e^{-M_1(\Delta t-t)} + A_0^* A_1 \langle 1|O_\Gamma|0\rangle e^{-M_1 t} e^{-M_0(\Delta t-t)} \end{aligned} \quad (3.1)$$

where Γ represents one of the sixteen Clifford elements defining the bilinears. To extract $\langle 0|O_\Gamma|0\rangle$ from the 2- and 3-point functions we examine five different fits. All errors are estimated with the full analysis done within a single-elimination jackknife method.

- *1-1 method* assumes a single state dominates the 2-pt and 3-pt functions. A_0 and M_0 are extracted from a fit to the 2-pt function and $\langle 0|O_\Gamma|0\rangle$ are estimated from the 3-pt functions keeping only the first term in Eq. (3.1).

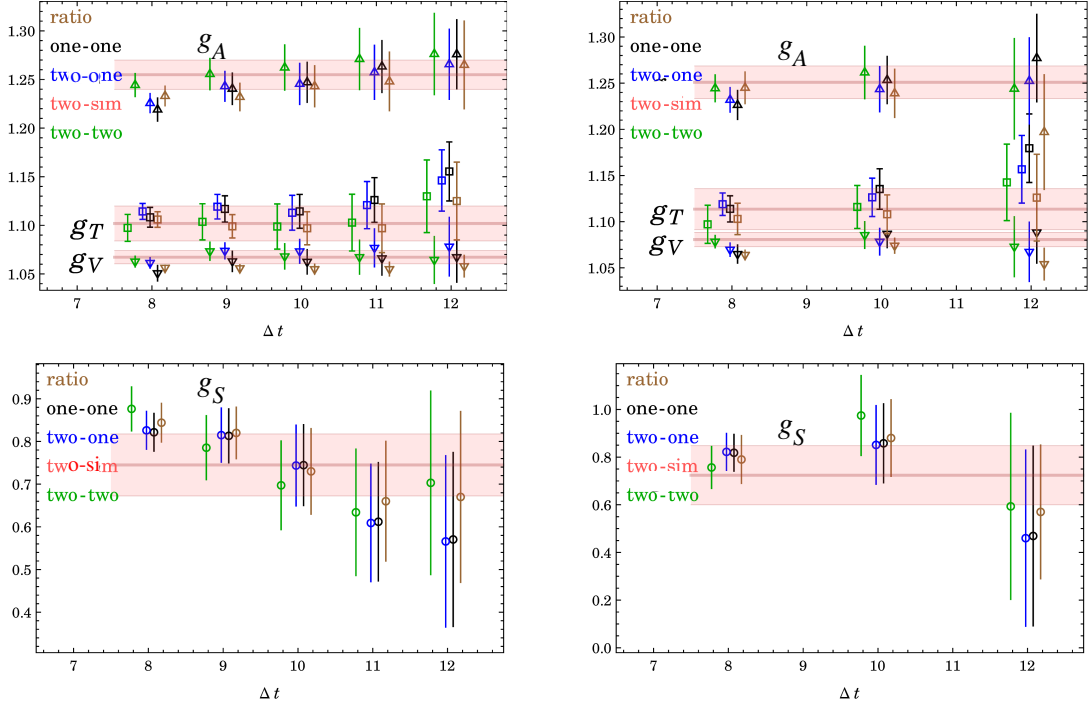


Figure 1: Estimates of the unrenormalized charges g_A , g_V , g_S and g_T from the five different fits for the $a = 0.12$ fm lattices with $M_\pi = 310$ MeV (left) and 220 MeV (right). The bands correspond to the 2-2 simultaneous fit described in the text and are used as our best estimates. The other four estimates are consistent with this estimate. Δt is the separation in time between the source and the sink in lattice units.

- *Ratio method* also assumes a single state dominates the 3-pt function. $\langle 0|O_\Gamma|0\rangle$ are estimated from the ratio of 3-pt to 2-pt functions. While some of the systematics cancel in the ratio, this fit relies on there being a good signal in the 2-pt function at separation Δt .
- *2-1 method*: A_0 , A_1 , M_0 and M_1 are extracted from a fit to the 2-pt function. Of these, only A_0 and M_0 are used to estimate $\langle 0|O_\Gamma|0\rangle$ by fitting to the first term in Eq. (3.1).
- *2-2 method* extracts A_0 , A_1 , M_0 and M_1 from a fit to the 2-pt function. These amplitudes and masses are used in a 2-parameter fit to the 3-pt function to estimate $\langle 0|O_\Gamma|0\rangle$ and $\langle 1|O_\Gamma|0\rangle$. We assume $\langle 0|O_\Gamma|1\rangle$ and $\langle 1|O_\Gamma|0\rangle$ are equal and we analyze only the real part of the 3-point function. We also neglect $\langle 1|O_\Gamma|1\rangle$.
- *2-2 simultaneous fit to all Δt* is the same as the 2-2 method for extracting A_0 , A_1 , M_0 and M_1 . The fit to the 3-point function uses data for all investigated values of Δt simultaneously.

Our current understanding of excited-state contamination is based on the analysis of the full set of 1013 configurations for the $a = 0.12$ fm $M_\pi = 310$ MeV ensemble using $\Delta t = 8, 9, 10, 11, 12$ and 958 configurations for the $a = 0.12$ fm $M_\pi = 220$ MeV ensemble using $\Delta t = 8, 10, 12$. These data are shown in Fig. 1. The vector charge g_V is shown only as a check.

We find that the ratio method exhibits plateaus for g_A , g_V and g_T , so it is not surprising that the first four methods give consistent estimates. The estimates from the 2-2 simultaneous fit method

are shown by the horizontal bands which we take as our best estimates. Our observations, based on these two ensembles of roughly 1000 configurations and smeared sources used by us in calculating quark propagators, are:

- On the 0.12 fm lattices, the statistical errors increase by about 40% with each unit increase in Δt . We find a similar increase for the $a = 0.09$ and 0.06 fm lattices once the unit increase in Δt is scaled by the factors 1.33 and 2 to keep distances constant in physical units.
- There is a small trend showing an increase (within 1σ) in g_A and g_T with Δt .
- For $t > 12$, even the signal in the 2-point nucleon correlator becomes noisy. Based on the trends seen in the $M_\pi = 310$ MeV ensemble, we considered it sufficient to investigate the $M_\pi = 220$ MeV ensemble using $\Delta t = 8, 10, 12$.
- The 2-2 *simultaneous fit* estimates of the central values and errors are consistent with data from other fits for all values of Δt .
- The errors increase by about 20% on lowering the light (u and d) quark masses by a factor of two *i.e.* from $M_\pi = 310$ to 220 MeV ensembles.
- The signal in g_S is the noisiest. On the $M_\pi = 220$ MeV ensembles, the error estimate is $\approx 20\%$, reasonably close to the desired accuracy.
- The statistical signal improves significantly as a is reduced. This is presumably due to the larger lattice volumes and smoother gauge configurations. Preliminary results show that on the $a = 0.09$ and 0.06 fm ensembles we will be able to extract even g_S within 10% accuracy.

Our conclusion, based on the $a = 0.12$ fm lattices, is that the central values from the 2-2 *simultaneous fit* agree with those from the other fits for separation $\Delta t = 10$ which corresponds to ≈ 1.2 fm. We, therefore, consider the $\Delta t = 10$ (*i.e.* in physical units this $\Delta t = 1.2$ fm) data the best compromise between reducing excited-state contamination and having a good statistical signal. The choice of Δt values investigated on the $a = 0.09$ and 0.06 fm ensembles is based on these conclusions.

4. Dependence on Quark Mass and Lattice Spacing

The data show a small increase in the values of the bare charges g_A , g_S and g_T on going from $M_\pi = 310$ to 220 MeV ensemble for $a = 0.12$ fm. The change is, however, within 1σ . The data for the $a = 0.09$ and 0.06 fm lattices are too preliminary to draw conclusions on the dependence on either the quark mass or on the lattice spacing. Since the statistical signal improves as $a \rightarrow 0$, we anticipate that our final results will shed light on possible dependence on the quark mass and the lattice spacing.

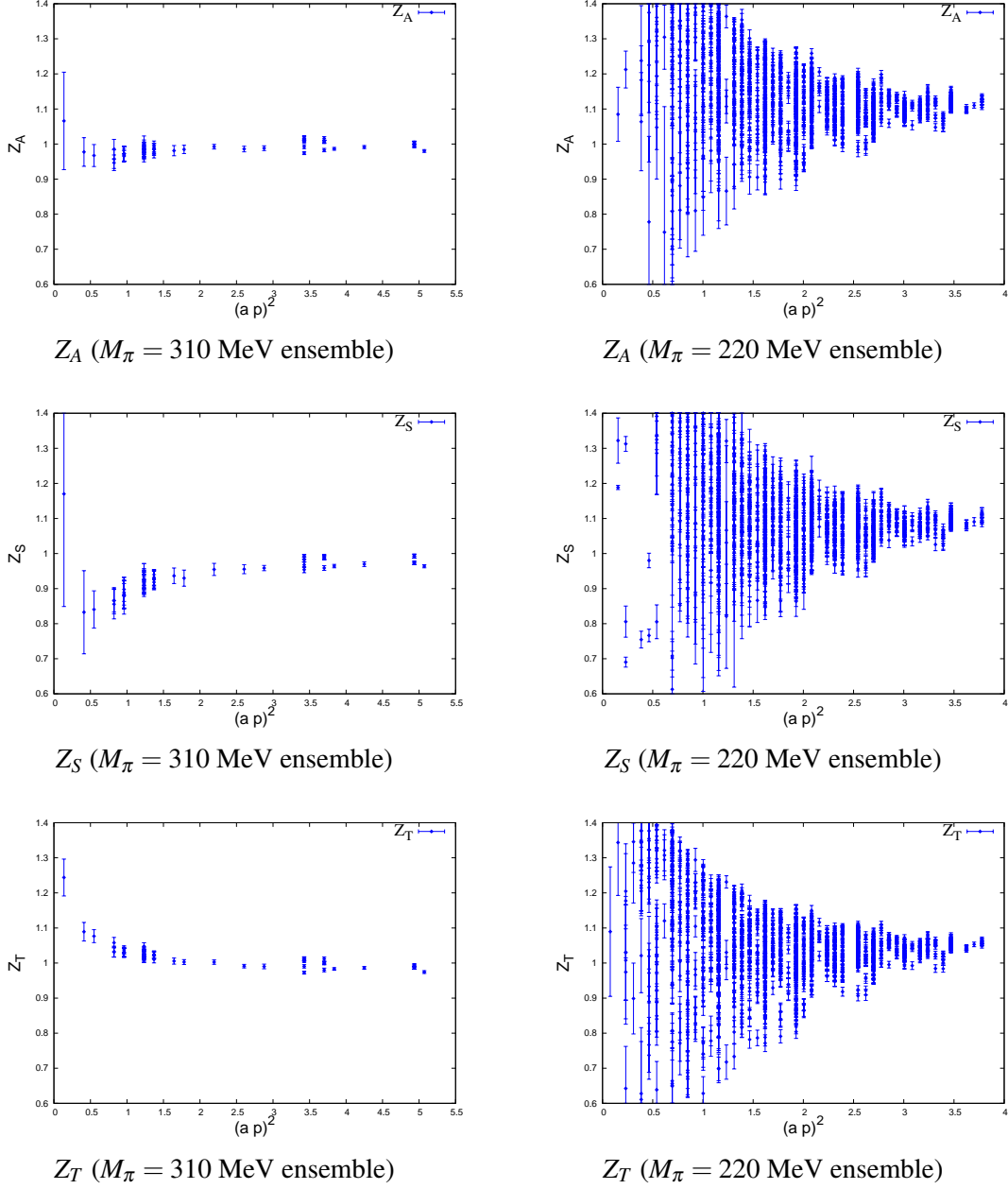


Figure 2: The data for the renormalization constants Z_A , Z_S and Z_T for the $a = 0.12$ fm lattices as a function of the momenta $p_i^2 = p_f^2 = (p_i - p_f)^2$. Note that $p = 1.6$ GeV for $pa = 1$. The analysis is done using point source Wilson-Clover action propagators calculated on HYP smeared HISQ lattices. For these ensembles, we find that estimates of all three Z_T are within 10% of unity. The dependence on the light quark mass, *i.e.*, the $\sim 10\%$ increase on going from the $M_\pi = 310$ to 220 MeV ensembles, is being studied.

5. Renormalization Constants Z_Γ

The renormalization constants Z_Γ are calculated in the RI-sMOM scheme using point source propagators evaluated on Landau gauge fixed lattices. The data for $a = 0.12$ fm ensembles are shown in Fig. 2. In this scheme only momenta satisfying $p_i^2 = p_f^2 = (p_i - p_f)^2$ are examined (the $M_\pi = 310$ MeV lattices have much fewer momenta satisfying this condition due to the smaller lattice volume), and the data have been averaged over points equivalent under cubic symmetry. We find that for the HYP smearing used in the calculation, all the renormalization constants are already within 10% of unity at $a = 0.12$ fm. We are currently investigating $O(p^4)$ contributions in these estimates to reduce these discretization effects.

6. Conclusions

We show that the excited-state contamination is smaller than statistical errors for source-sink separation greater than 1.2 fm. The renormalization constants in the RI-sMOM scheme are within 10% of unity already on the coarsest lattices at $a = 0.12$ fm. We therefore, conclude that $O(1000)$ lattices will be sufficient to get estimates with 10–15% precision for both g_S and g_T at all three values of the lattice spacing. We anticipate completing the analysis on the $a = 0.09$ and 0.06 fm lattices over the next year. These data are needed to confirm the small observed dependence on the quark mass and lattice spacing, and before conclusions on the chiral and continuum extrapolations can be drawn.

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